

Noether's Second Theorem and Strange Metals

Thanks to: NSF

Gabriele La Nave

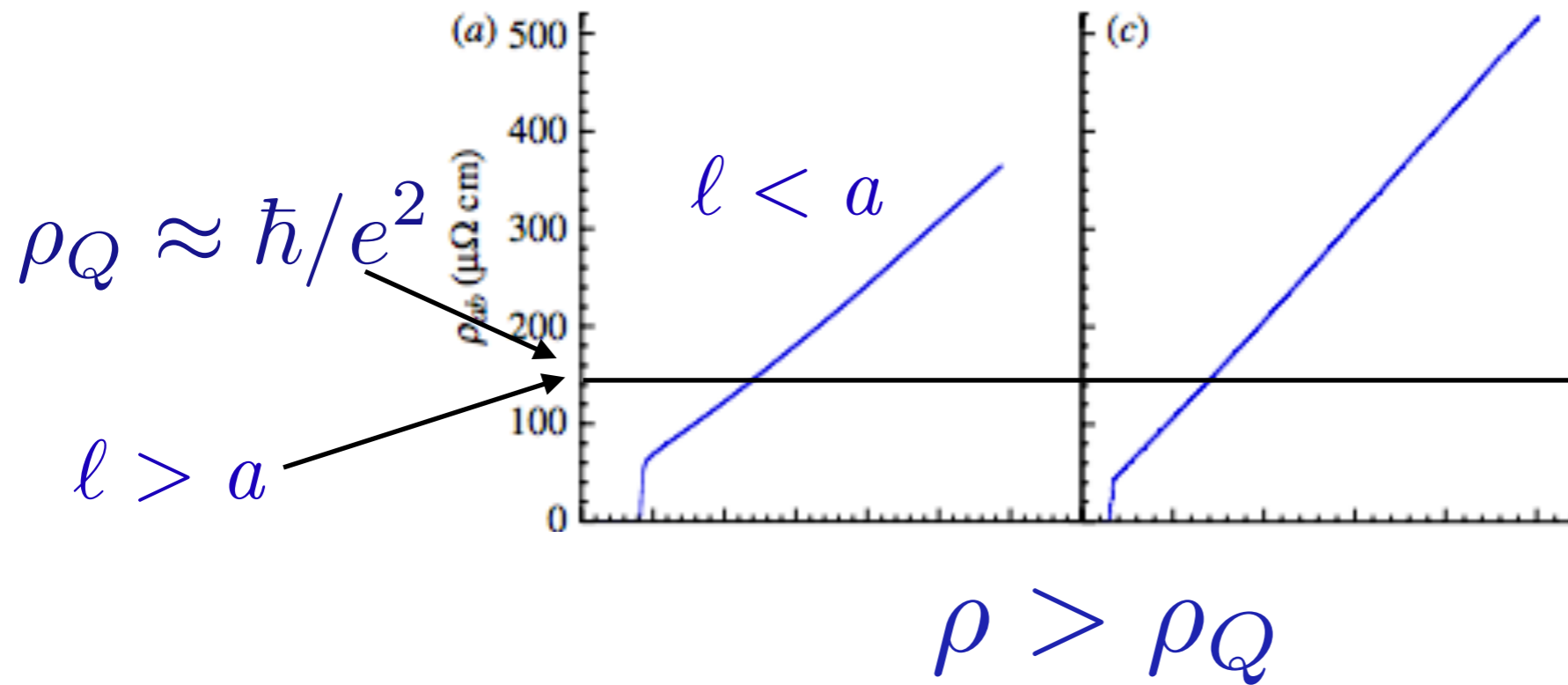


Rev. Mod. Phys. 2019
(arXiv:1904.01023)
CIMP 2019
Adv. Th. Math. Phys.
2019



Kridsangaphong Limtragool

T-linear resistivity



Violates MIR limit

electrons not charge carriers

?

Does anything
local carry the charge?

No

$$\int_{\ell} eA \notin h\mathbb{Z}$$

new gauge
principle?

Yes

$$\oint eA \cdot dl \in h\mathbb{Z}$$

$$e = ?$$

Can you mess with
A?

Noether's First Theorem

$$U(1) \longleftrightarrow qA \rightarrow qA - q\partial_\mu\Lambda \quad \psi' = e^{iq\Lambda}\psi$$

$$[q\Lambda] = 0$$

$$[qA] = 1$$

$$[A] = 1$$

$$S = \int d^d x (J_\mu A^\mu + \dots)$$

fixes dimension of current

$$S \rightarrow S + \int d^d x J_\mu \partial\Lambda$$

$$[d^d x J A] = 0$$

$$[J] = d - 1$$

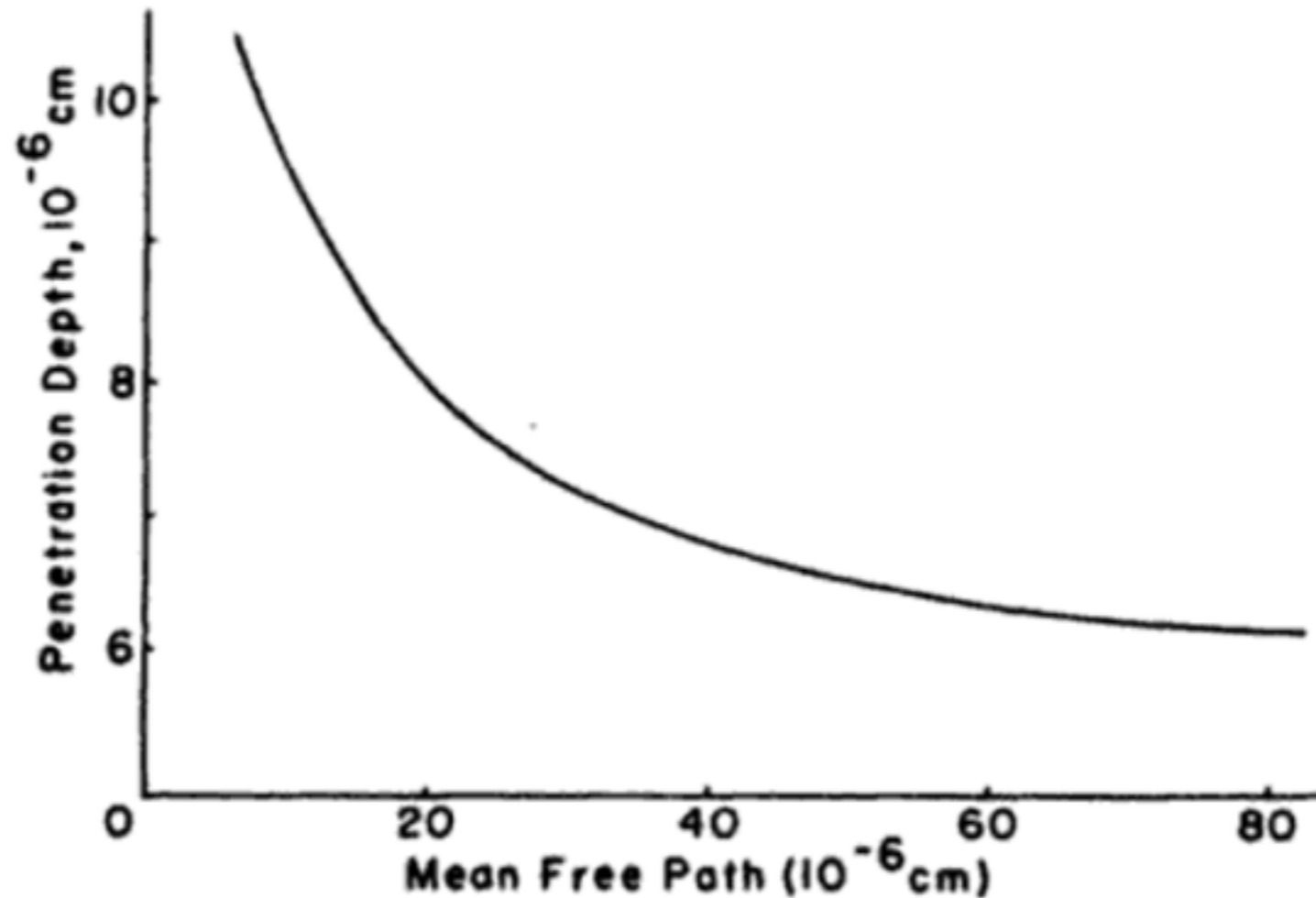
$$\partial_\mu J^\mu = 0$$

Noether's Thm. I

current conservation

Are there exceptions?

Pippard's problem



$$J_s \neq \frac{-c}{4\pi\lambda^2} A$$

London Eq.

failure of local London relations

Superconductivity ala Weinberg

$$U(1) \rightarrow \mathbb{Z}_2$$

$$A_\mu - \partial_\mu \phi = 0$$

$U(1)/\mathbb{Z}_2$

$$\nabla \phi - A = 0$$

stable equilibrium

around
minimum

$$L_m = L_{m0} - \frac{1}{2} \int C^{\mu\nu}(\mathbf{x}, \mathbf{x}') (A_\mu(\mathbf{x}) - \partial_\mu \phi(\mathbf{x})) \\ \times (A_\nu(\mathbf{x}') - \partial_\nu \phi(\mathbf{x}')) d^3 \mathbf{x}' d^3 \mathbf{x} + \dots$$

Pippard Current

$$J_{\mu}(\mathbf{x}) = \frac{\delta \mathbf{L}_m}{\delta \mathbf{A}_{\mu}} = - \int \mathbf{C}^{\mu\nu}(\mathbf{x}, \mathbf{x}') (\mathbf{A}_{\nu}(\mathbf{x}') - \partial_{\nu} \phi(\mathbf{x}')) d^3 \mathbf{x}'$$



Pippard
kernel

$$J_s = - \frac{3}{4\pi c \xi_0 \lambda} \int \frac{(\vec{r} - \vec{r}') ((\vec{r} - \vec{r}') \cdot \vec{A}(\vec{r}')) e^{-(\vec{r} - \vec{r}')/\xi(\ell)}}{(\vec{r} - \vec{r}')^4} d^3 \vec{r}'$$

non-local

Units of Current

$$J_\mu(\mathbf{x}) = \frac{\delta \mathbf{L}_m}{\delta \mathbf{A}_\mu} = - \int \mathbf{C}^{\mu\nu}(\mathbf{x}, \mathbf{x}') (\mathbf{A}_\nu(\mathbf{x}') - \partial_\nu \phi(\mathbf{x}')) d^3 \mathbf{x}'$$

$$[J] = d - d_C - d_A$$

anomalous
dimension

Standard Result

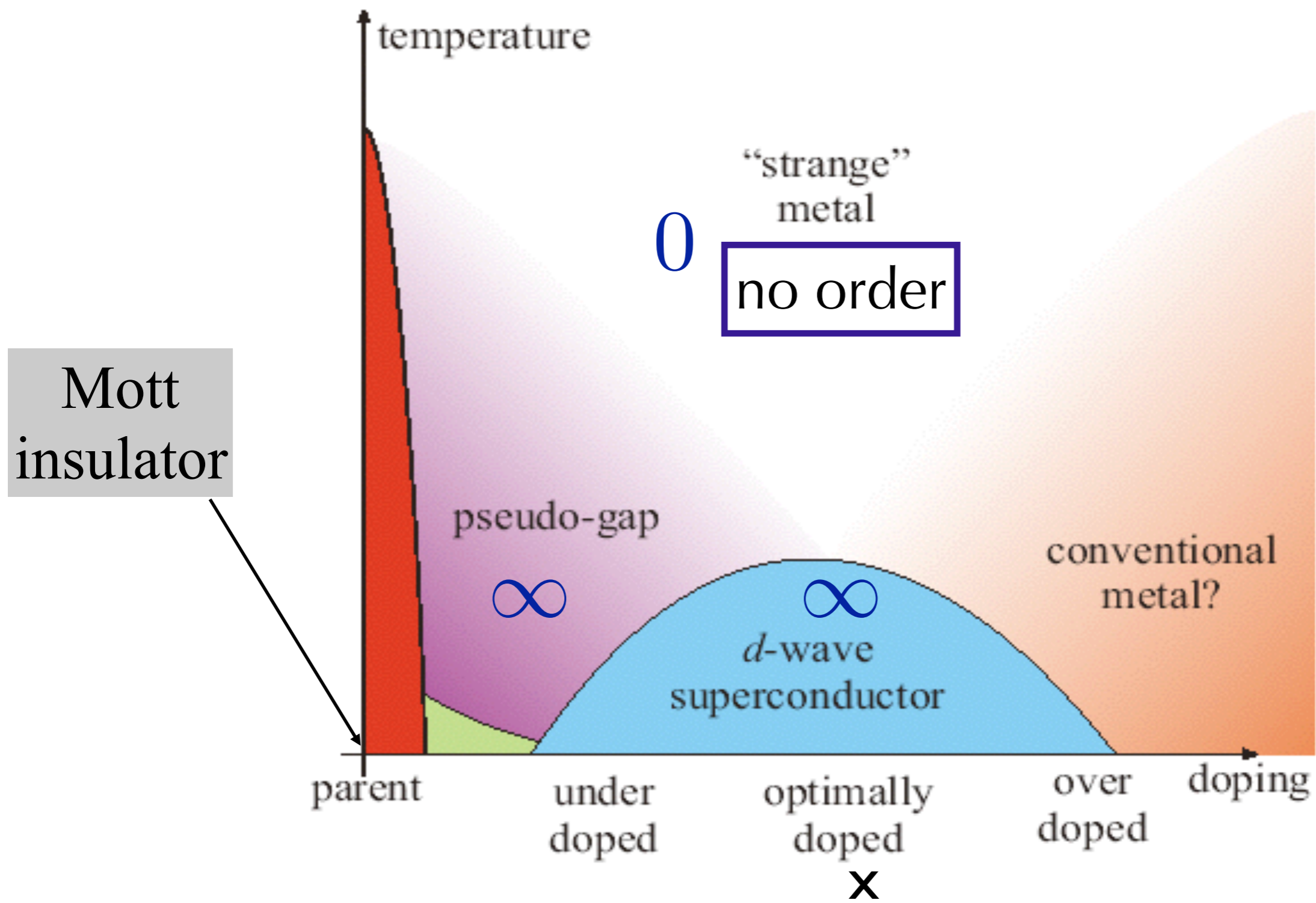
$$\delta(x_0 - y_0) [J_\mu(x), \phi(y)] = \delta^d(x - y) \delta\phi(y)$$

$$[J] = d - 1$$

Are there other
examples of
currents with
anomalous
dimensions?

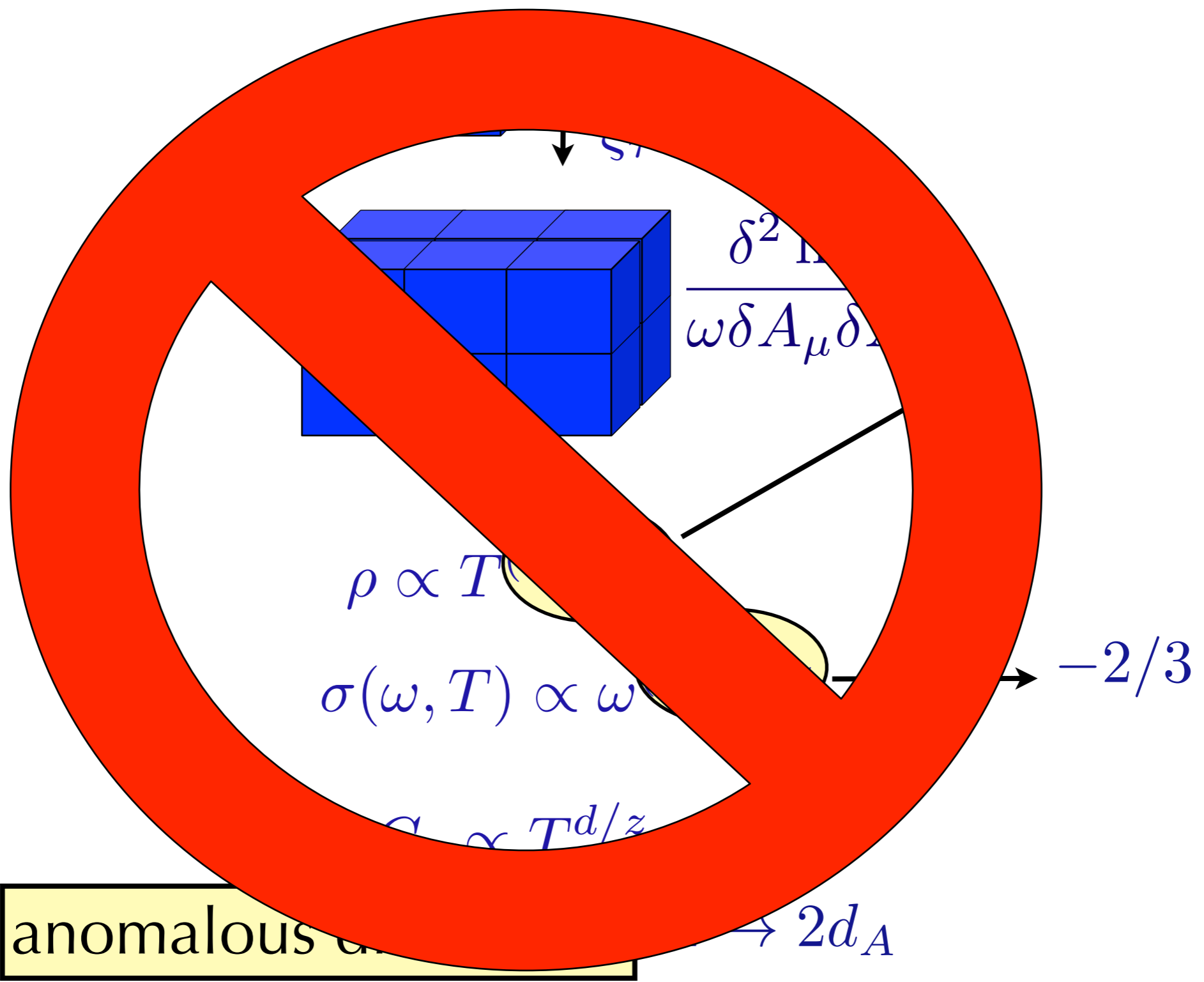
underlying
electricity and
magnetism?

is symmetry
breaking
necessary?



why is the problem hard?

single-parameter scaling

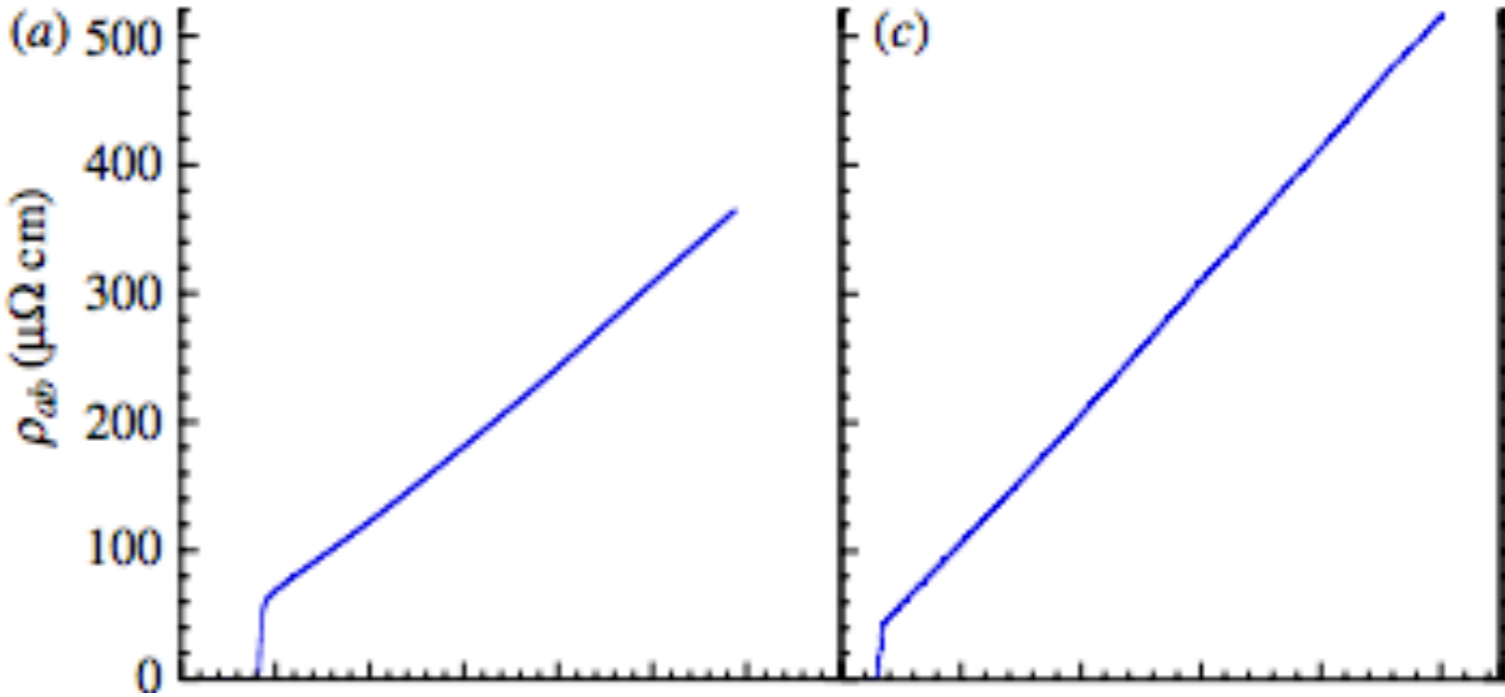


strange metal explained!

Hall Angle

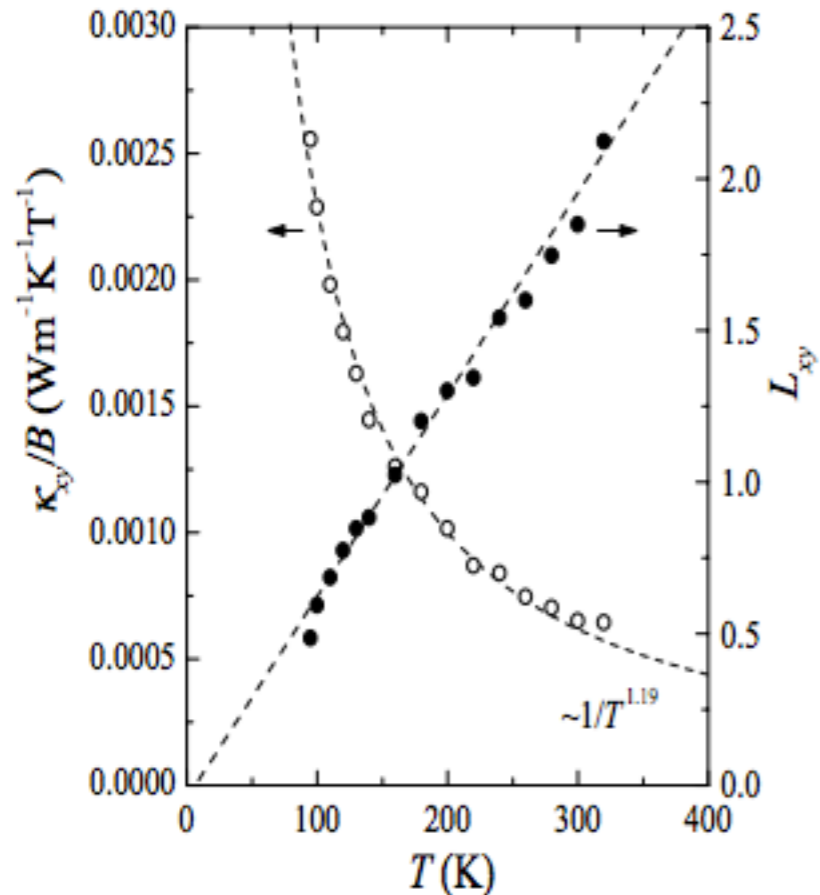
$$\cot \theta_H \equiv \frac{\sigma_{xx}}{\sigma_{xy}} \approx T^2$$

T-linear resistivity



Hall Lorenz ratio

$$L_{xy} = \kappa_{xy} / T \sigma_{xy} \neq \# \propto T$$



all explained if

$$[J_\mu] = d - \theta + \Phi + z - 1$$

Hartnoll/Karch

$$[A_\mu] = 1 - \Phi$$

$$\Phi = -2/3$$

strange metal

$$[J_\mu] = d - \theta + \Phi + z - 1$$

$$\begin{aligned} [A_\mu] &= 1 - \Phi \\ \Phi &= -2/3 \\ [E] &= 1 + z - \Phi \\ [B] &= 2 - \Phi \end{aligned}$$



note $\pi r^2 B \neq \text{flux}$

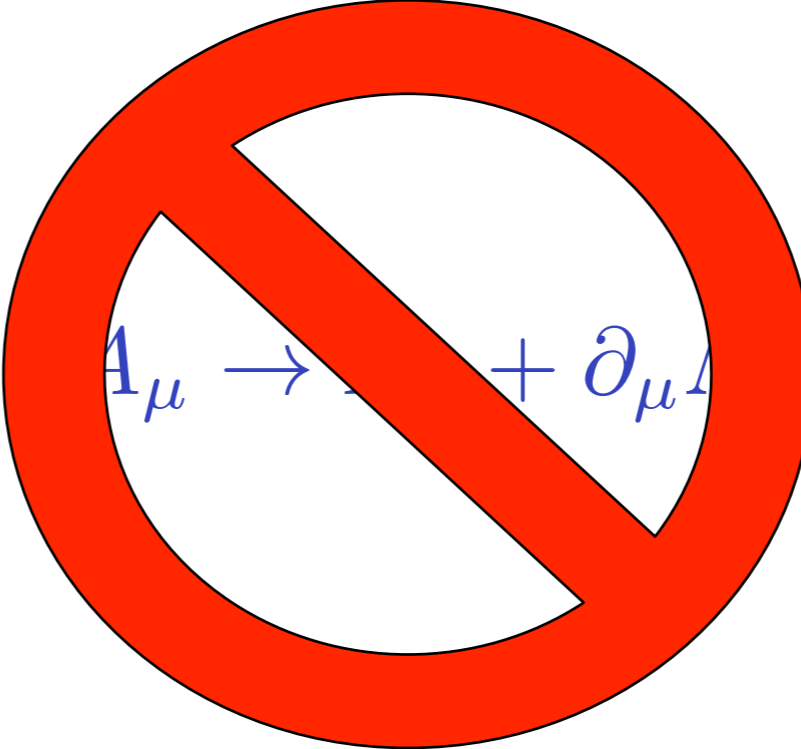
$$\oint A \cdot d\ell \notin h\mathbb{Z}$$

How is this
possible - -
if at all?

what is the new gauge principle?

if

$$[A_\mu] \neq 1$$


$$A_\mu \rightarrow A_\mu + \partial_\mu \Lambda$$

Noether's Second Theorem: precursor

hint

$$\partial_\mu J^\mu = 0$$

current conservation

what if

$$[\partial_\mu, \hat{Y}] = 0$$

new current

gauge symmetry

$$\partial_\mu \hat{Y} J^\mu = \partial_\mu \tilde{J}^\mu = 0$$

$$[\tilde{J}] = d - 1 - D_Y$$

possible gauge transformations

$$S = -\frac{1}{4} \int d^d x F^2$$



$$S = \frac{1}{2} \int \frac{d^d k}{2\pi^d} A_\mu(k) \underbrace{[k^2 \eta^{\mu\nu} - k^\mu k^\nu]}_{M^{\mu\nu} k_\nu = 0} A_\nu(k)$$

$$M^{\mu\nu} k_\nu = 0$$

zero eigenvector

$$ik_\mu \rightarrow \partial_\nu$$

$$A_\mu \rightarrow A_\mu + \partial_\mu \Lambda$$

family of zero eigenvalues

$$M_{\mu\nu} \underbrace{f k^\nu} = 0$$

generator of gauge symmetry

- 1.) rotational invariance
- 2.) A is still a 1-form
- 3.) $[f, k_\mu] = 0$

only choice

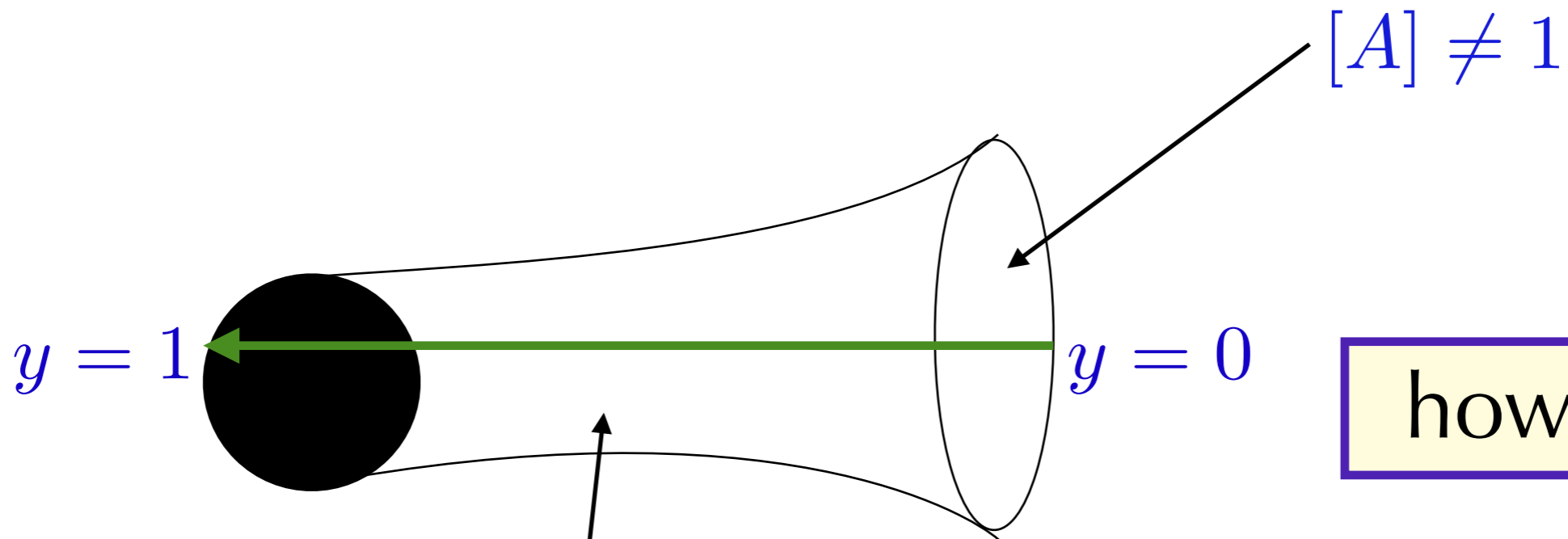
$$f \equiv f(k^2)$$

$$(-\Delta)^\gamma$$

$$A_\mu \rightarrow A_\mu + (-\Delta)^{\frac{\gamma-1}{2}} \partial_\mu \Lambda \quad [A_\mu] = \gamma$$

what kind of E&M has such
gauge transformations?

claim: extra dimension



how?



$$S = \int dV_d dy (y^a F^2 + \dots)$$

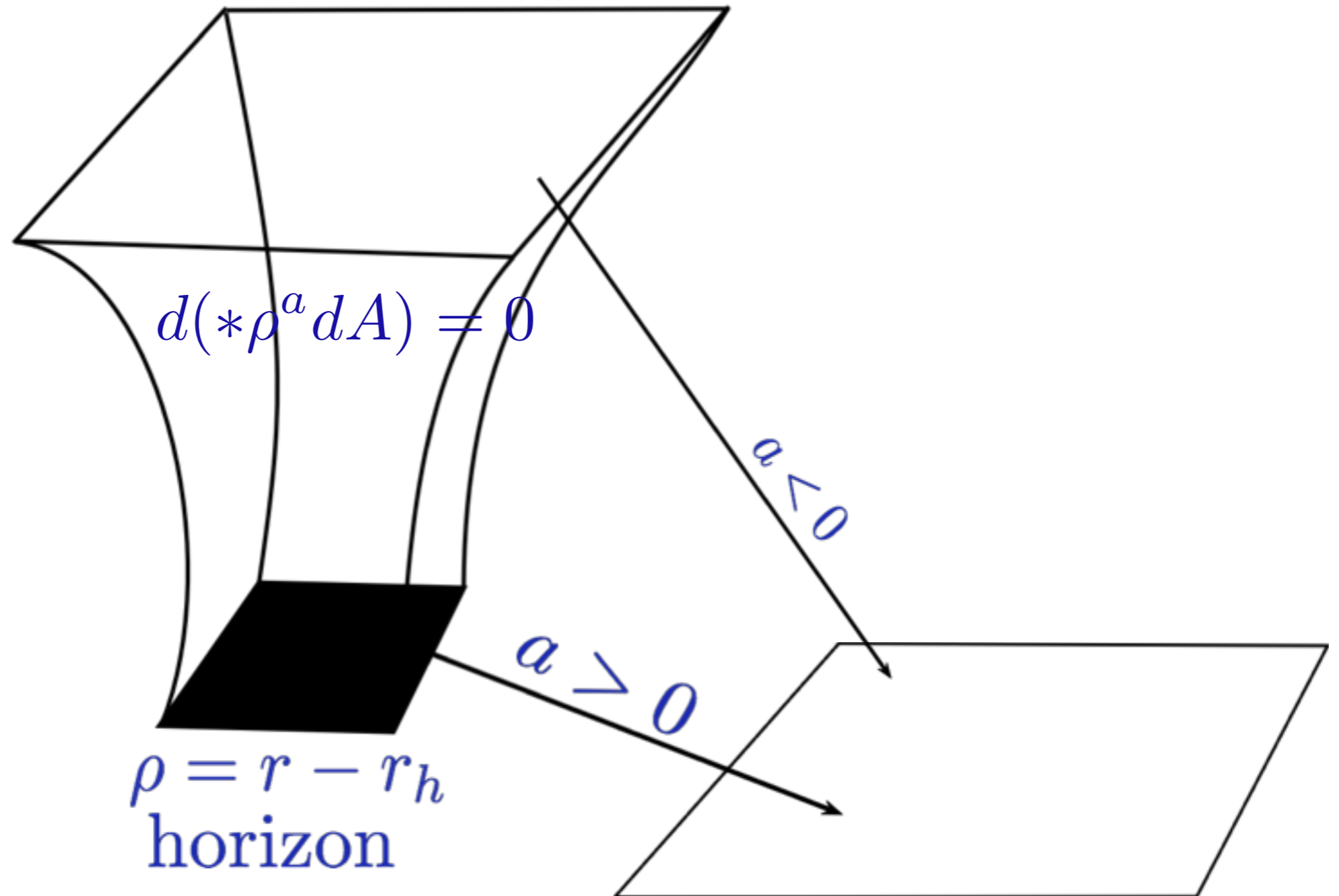
eom $d(y^a \star dA) = 0$

Karch:1405.2926
Gouteraux: 1308.2084

if holography is RG then
how can it lead to an
anomalous dimension?

membrane paradigm

conformal boundary
 $r \rightarrow \infty$



construct 'boundary'
theory explicitly

Caffarelli-Silvestre
extension theorem
(2006)

y

$$g(x, y = 0) = f(x)$$

$$\Delta_x g + \frac{a}{y} g_y + g_{yy} = 0$$

$$\nabla \cdot (y^a \nabla g(x, y)) = 0$$

$$\lim_{y \rightarrow 0} y^a \partial_y g$$

?

$$C_{d,\gamma} (-\Delta)^\gamma f$$

x

fractional Laplacian

$$g(z = 0, x) = f(x)$$

$$\gamma = \frac{1 - a}{2}$$

Caffarelli-Silvestre
extension theorem
(2006)

Dirichlet

\mathbb{R}^{n+1} $g(x, y)$



$(-\Delta)^\gamma$

\mathbb{R}^n $f(x)$

Neumann

closer look

$$\nabla \cdot (y^a \nabla u) = 0$$

scalar field
(use CS theorem)

$$d(y^a \star dA) = 0$$

holography

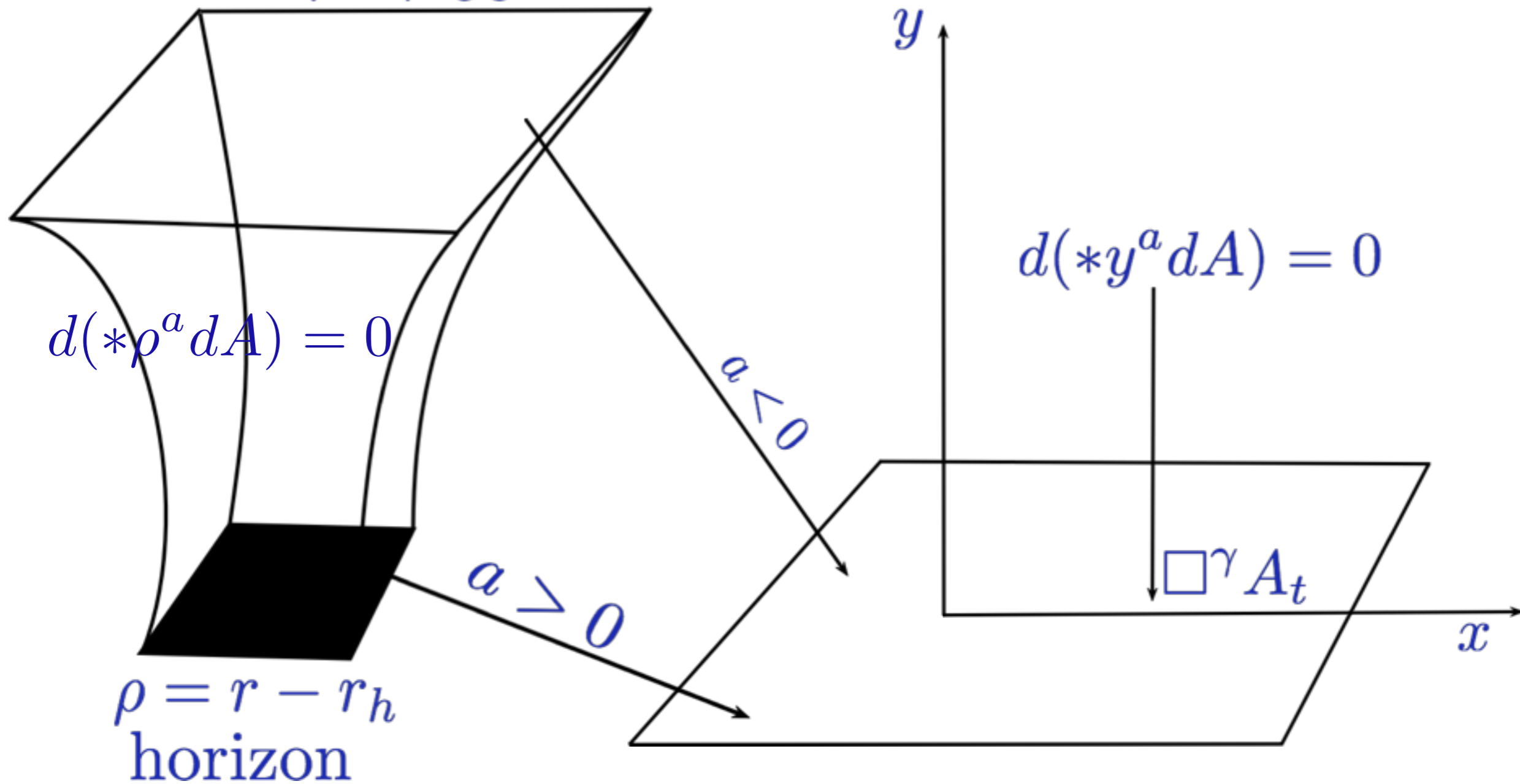
similar equations

generalize CS theorem to
p-forms

GL,PP:1708.00863

(CIMP, 366, 199 (2019))

UV
conformal boundary
 $r \rightarrow \infty$



$\rho = r - r_h$
horizon

IR

$$A \rightarrow A + d_\gamma \Lambda \equiv A'$$

$$d_\gamma \equiv \square^{\frac{\gamma-1}{2}} d$$

boundary action:
fractional Maxwell
equations

$$\square^\gamma A_\perp = J$$

boundary action has
'anomalous dimension'
(non-locality)

$$F \rightarrow d_\gamma A = \partial_\mu \square^{(\gamma-1)/2} A_\nu - \partial_\nu \square^{(\gamma-1)/2} A_\mu$$

if holography is RG then
how can it lead to an
anomalous dimension?

$$S = \int dV_d dy (y^a F^2 + \dots)$$



$$[A] = 1 - a/2$$

dimension of A is fixed by
the bulk theory: not really
anomalous dimension

new gauge transformation

$$A \rightarrow A + d_\gamma \Lambda \equiv A'$$

$$d_\gamma \equiv \square^{\frac{\gamma-1}{2}} d$$

$$[A] = \gamma$$

action of gauge group

$$D_{\gamma,A}(e^\Lambda \odot \phi) = e^{i\square^{(1-\gamma)/2}\Lambda} D_{\gamma,A'}\phi$$

$$D_{\gamma,A}\phi = \left(\partial_\mu + ie\square^{(1-\gamma)/2} A^\mu \right) \square^{(1-\gamma)/2} \phi$$

$$L = D_{\gamma,A}\phi(D_{\gamma,A}\phi)^* - m^2\phi^*\phi - F_{\gamma}^{\mu\nu}F_{\mu\nu\gamma}$$

from the bulk

use CS theorem

$$\lim_{y \rightarrow 0} [y^a \partial \phi(x, y), y^a \partial \phi(x', y)] = \Delta^\gamma [\phi(x, y), \phi(x', y)] \\ = \Delta^\gamma \delta(x - x') = 0$$

no problem
with causality

Ward identities

$$C^{ij}(k) \propto (k^2)^\gamma \left(\eta^{ij} - \frac{k^i k^j}{k^2} \right).$$

standard Ward
identity

$$k_i C^{ij}(k) = 0 \quad \longrightarrow \quad \partial_i C^{ij}(k) = 0$$

but

$$k^{\gamma-1} k_\mu C^{\mu\nu} = 0 \quad \longrightarrow \quad \partial_\mu (-\Delta)^{\frac{\gamma-1}{2}} C^{\mu\nu} = 0$$

inherent ambiguity in E&M

Noether's Second Theorem

$$\begin{aligned}
 & \sum \psi_{\mathbf{i}} \delta u_{\mathbf{i}} = \delta \mathcal{F} - \\
 & - \frac{d}{dx} \left\{ \sum \left[\binom{1}{1} \frac{\partial \mathcal{F}}{\partial u_{\mathbf{i}}^{(1)}} \delta u_{\mathbf{i}} + \binom{2}{1} \frac{\partial \mathcal{F}}{\partial u_{\mathbf{i}}^{(2)}} \delta u_{\mathbf{i}}^{(1)} + \dots + \binom{\kappa}{1} \frac{\partial \mathcal{F}}{\partial u_{\mathbf{i}}^{(\kappa)}} \delta u_{\mathbf{i}}^{(\kappa-1)} \right] \right\} + \\
 & + \frac{d^2}{dx^2} \left\{ \sum \left[\binom{2}{2} \frac{\partial \mathcal{F}}{\partial u_{\mathbf{i}}^{(2)}} \delta u_{\mathbf{i}} + \binom{3}{2} \frac{\partial \mathcal{F}}{\partial u_{\mathbf{i}}^{(3)}} \delta u_{\mathbf{i}}^{(1)} + \dots + \binom{\kappa}{2} \frac{\partial \mathcal{F}}{\partial u_{\mathbf{i}}^{(\kappa)}} \delta u_{\mathbf{i}}^{(\kappa-2)} \right] \right\} + \\
 & \vdots \\
 & + (-1)^\kappa \frac{d^\kappa}{dx^\kappa} \left\{ \sum \left[\binom{\kappa}{\kappa} \frac{\partial \mathcal{F}}{\partial u_{\mathbf{i}}^{(\kappa)}} \delta u_{\mathbf{i}} \right] \right\}
 \end{aligned} \tag{6}$$

$$A_\mu \rightarrow A_\mu + \partial_\mu \Lambda + \partial_\mu \partial_\nu G^\nu + \dots,$$

$$A \rightarrow A + d_\gamma \Lambda \equiv A'$$

$$d_\gamma \equiv (-\Delta)^{\frac{\gamma-1}{2}} d$$

$$\Delta \rightarrow \square$$

Noether's Second Theorem and Ward Identities for Gauge Symmetries

Steven G. Avery^a, Burkhard U. W. Schwab^b

For simplicity, we focus on the case when the transformation may be written in the form⁶

$$\delta_\lambda \phi = f(\phi) \lambda + f^\mu(\phi) \partial_\mu \lambda, \quad (10)$$

but it is straightforward to consider transformations, as Noether did, involving arbitrarily high derivatives of λ . (Although, the authors know of no physically interesting examples.) Let us start with

arxiv:1510.07038

family of zero eigenvalues

$$M_{\mu\nu} f k^\nu = 0$$

most fundamental conservation law

$$\partial^\mu \underbrace{(-\nabla^2)^{(\gamma-1)/2} J_\mu}_{J'_\mu} = 0$$

is boundary non-locality a problem?

entanglement?

Li/Takayanagi (PRL 106, 141301 (2011))

$$I_B(\phi) = \int d^d x \phi (-\Delta)^\gamma \phi$$



$$S_B = \kappa_{d-2} \left(\frac{1}{\epsilon} \right)^{d-2}$$

Area

$$I_C(\phi) = \int d^d x \phi e^{(-\Delta)^\gamma} \phi$$

?



$$S_C \sim \kappa_{d-2} \left(\frac{1}{\epsilon} \right)^{d-2+2\gamma}$$

>Area

RESOLUTION

$$I = \int d^d x \phi \mathcal{O}^\gamma \phi + J \phi$$

$$\phi \rightarrow \mathcal{O}^{(1-\gamma)/2} \phi$$

$$J \rightarrow \mathcal{O}^{(1-\gamma)/2} J$$

Rule

$$\mathcal{O}(\hat{\gamma}) = \hat{\mathcal{O}}^\gamma$$

$$I = \int d^d x \phi \mathcal{O} \phi + J \phi$$

Area
entanglement

simplest exception

standard AdS

$$I = \int d^d x \phi \mathcal{O}^\gamma \phi + m^2 |\phi|^2$$

$$m^2 > 0$$

$$S = \kappa_{d-2} \left(\frac{1}{\epsilon} \right)^{d-2\gamma}$$

Volume
entanglement
 $\gamma = 1/2$

is there a consistent algebra
for fractional currents?

Yes

Virasoro algebra

$$L_n := -z^{n+1} \frac{\partial}{\partial z}$$

$$[L_n, L_m] = (n - m)L_{n+m} + \frac{c}{12}m(m^2 - 1)\delta_{m+n,0}$$

Witt algebra

central
extension

conformal
transformations
on unit disk

$$\mathcal{V} \rightarrow \mathcal{W} \rightarrow 1$$

Fractional Virasoro algebra

generators

$$L_n^a = -z^{a(n+1)} \left(\frac{\partial}{\partial z} \right)^a \quad \bar{L}_n^a := -\bar{z}^{a(n+1)} \left(\frac{\partial}{\partial \bar{z}} \right)^a$$

$$[L_n, L_m](z^{ak}) = \left(\frac{\Gamma(a(k+n)+1)}{\Gamma(a(k-1+n)+1)} - \frac{\Gamma(a(k+m)+1)}{\Gamma(a(k-1+m)+1)} \right) L_{n+m}(z^{ak})$$

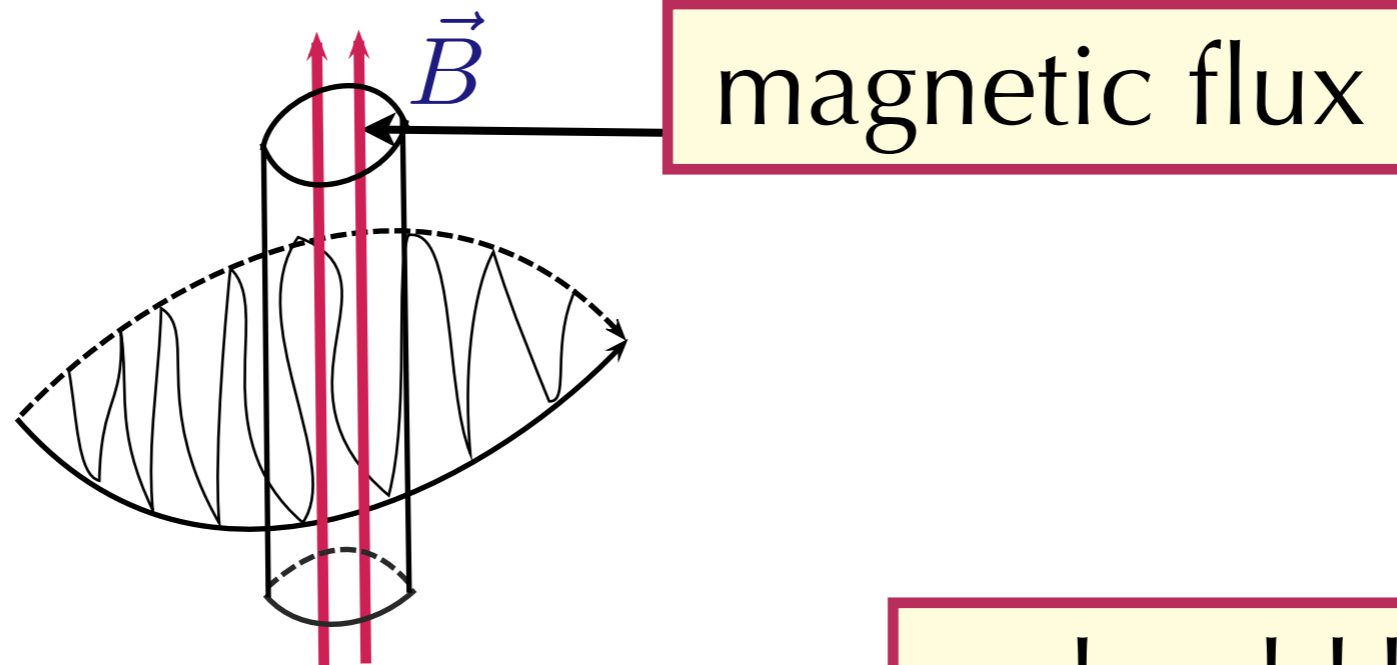
$$= (A_{n,m}^a(k) \otimes L_{n+m})(z^{ak})$$

$$[L_m^a, L_n^a] = A_{m,n} L_{m+n}^a + \delta_{m,n} h(n) c Z^a$$

algebra for conformal non-local actions

$$Z_\star^2(\mathcal{W}_a, \mathcal{H}) / B_\star^2(\mathcal{W}_a, \mathcal{H})$$

experiments?



magnetic flux

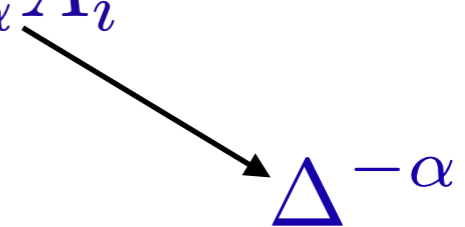
$$\pi r^2 B$$

should be dimensionless

$$[B] = 2 - \Phi = 2 + 2/3 \neq 2$$

what's the resolution?

correct dimensionless quantity

$$a_i \equiv [\partial_i, I_i^\alpha A_i] = \partial_i I_i^\alpha A_i$$


what's the relationship?

$$\oint_{\partial\Sigma} a \qquad \oint_{\partial\Sigma} A$$

$$\text{Norm} \oint_{\partial\Sigma} a = \frac{1}{\Gamma(3/2 - \gamma)} \oint_{\partial\Sigma} A$$

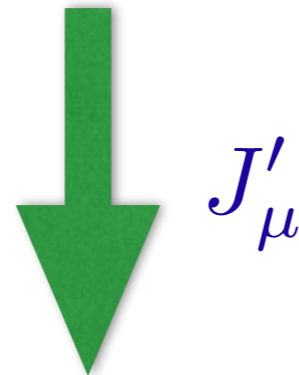
not an
integer

obstruction theorem to charge quantization (NST)

$$A_\mu \rightarrow A_\mu + \partial_\mu \Lambda + \partial_\mu \partial_\nu G^\nu + \dots,$$

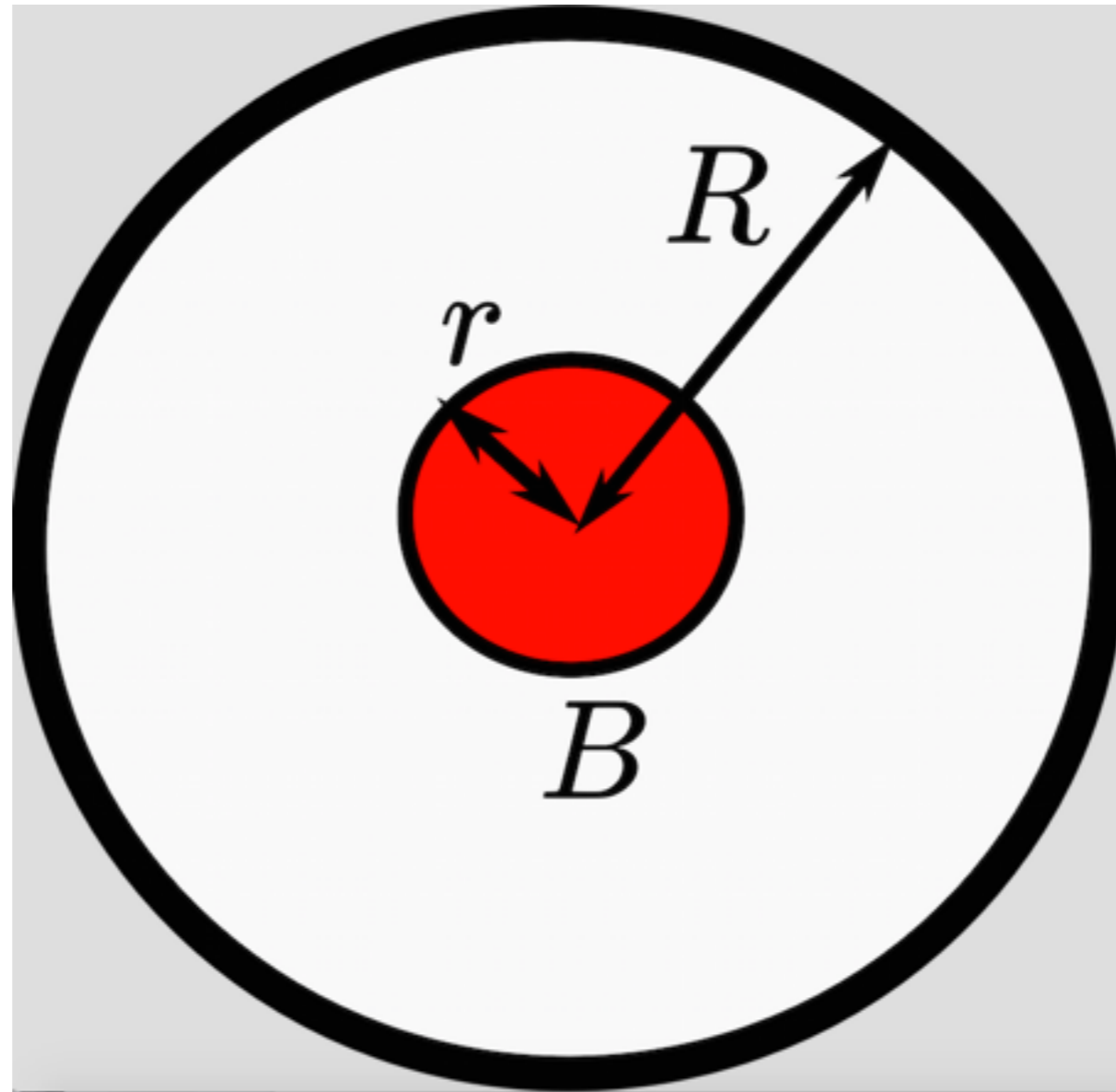
$$A \rightarrow A + d_\gamma \Lambda \equiv A'$$

$$d_\gamma \equiv (-\Delta)^{\frac{\gamma-1}{2}} d$$



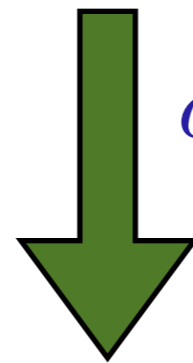
charge ill-defined (new
landscape problem)

New Aharonov-Bohm Effect



$$\Delta\phi_D = \frac{e}{\hbar} \pi r^2 B R^{2\alpha-2} \left(\frac{\sqrt{\pi} 2^{1-\alpha} \Gamma(2-\alpha) \Gamma(1-\frac{\alpha}{2})}{\Gamma(\alpha) \Gamma(\frac{3}{2}-\frac{\alpha}{2})} \sin^2 \frac{\pi\alpha}{2} {}_2F_1(1-\alpha, 2-\alpha; 2; \frac{r^2}{R^2}) \right)$$

is the correction large?



$$\alpha = 1 + 2/3 = 5/3$$

$$\Delta\Phi_R = \frac{eB\ell^2}{\hbar} L^{-5/3} / (0.43)^2$$

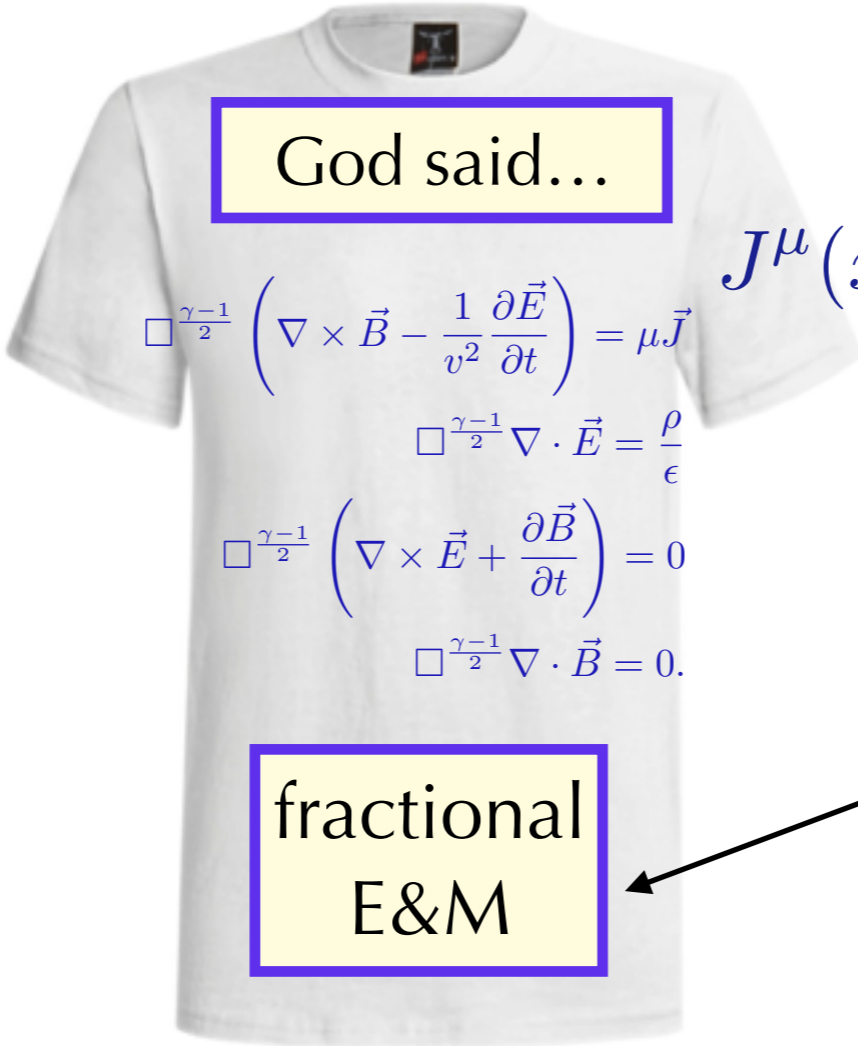
yes!

if in the strange metal



$$[A_\mu] = d_A \neq 1$$

Pippard Kernel



God said...

$$\square^{\frac{\gamma-1}{2}} \left(\nabla \times \vec{B} - \frac{1}{v^2} \frac{\partial \vec{E}}{\partial t} \right) = \mu \vec{J}$$

$$\square^{\frac{\gamma-1}{2}} \nabla \cdot \vec{E} = \frac{\rho}{\epsilon}$$

$$\square^{\frac{\gamma-1}{2}} \left(\nabla \times \vec{E} + \frac{\partial \vec{B}}{\partial t} \right) = 0$$

$$\square^{\frac{\gamma-1}{2}} \nabla \cdot \vec{B} = 0.$$

$$J^\mu(x) = - \int d^d x' C_{\mu\nu}(|x - x'|) A^\nu$$

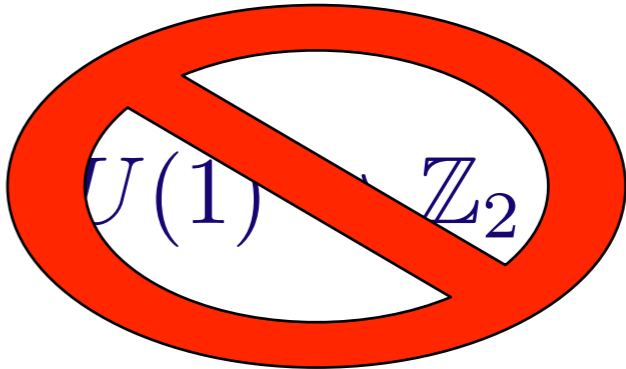
$$[J] \neq d - 1$$

$$[A] \neq 1$$

fractional E&M

in SC!

$$\omega = ck$$



$U(1) \rightarrow \mathbb{Z}_2$